

# The Cahn-Hilliard equation in image inpainting

Alain Miranville

Université de Poitiers, France

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The Cahn–Hilliard equation : recent advances and applications

Applications of the Cahn-Hilliard equation in image processing :

- Image denoising
- Image inpainting





## **A Cahn-Hilliard model for binary image inpainting (A. Bertozzi-S. Esedoglu-A. Gillette, 2007)**

Image inpainting : consists in filling in parts of an image/video from the surrounding area (interpolation)

Applications : restoration of old paintings, removing scratches, altering scenes, restoration of motion pictures, ...

PDE's for image inpainting : M. Bertalmio et al. (Navier-Stokes like model)

Other models : second-order models (S. Esedoglu-J. Shen)



$$g(x, s) = \lambda_0(s - h(x))\chi_{\Omega \setminus D}(x), \quad \lambda_0 > 0, \quad D \subset\subset \Omega$$

Equation :

$$\frac{\partial u}{\partial t} + \epsilon \Delta^2 u - \frac{1}{\epsilon} \Delta f(u) + g(x, u) = 0, \quad \epsilon > 0$$

$h(x)$  : given image ( $h \in L^2(\Omega)$ )

$D$  : inpainting domain (damaged region)

$$f(s) = 4s^3 - 6s^2 + 2s$$

$g(x, u)$  : added to keep  $u$  close to the image  $h(x)$  outside the inpainting region (fidelity term)

Advantages (over, e.g.,  $u = h$  outside  $D$ ) : no regularity assumption on  $D$ , no perfect  $h$  outside  $D$





Algorithm : dynamic two-steps algorithm involving  $\epsilon$

First step : large value of  $\epsilon$  to connect the edges

Second step : small value of  $\epsilon$  (depending on the mesh size); solution obtained in the first step as initial datum

Idea : solve the equation to steady state to construct an inpainting version  $u(x)$  of  $h(x)$

Advantage : fast in numerical simulations

**Remark :** Limit problem when  $\lambda_0 = +\infty$  ( $g$  of class  $\mathcal{C}^2$ ) :

$$\Delta(\epsilon \Delta u - \frac{1}{\epsilon} f(u)) = 0 \text{ in } D$$

$$u = h \text{ on } \partial D$$

$$\nabla u = \nabla h \text{ on } \partial D$$

→ Continuation of the image gradient into the missing domain

→  $\lambda_0$  large in numerical simulations

**Mathematical analysis** (L. Cherfils-H. Fakhir-A. Miranville, IPI, SIIMS) :

$$\begin{aligned}\frac{\partial u}{\partial t} + \Delta^2 u - \Delta f(u) + \chi_{\Omega \setminus D}(x)u &= 0 \\ \frac{\partial u}{\partial \nu} &= \frac{\partial \Delta u}{\partial \nu} = 0 \text{ on } \Gamma \\ u|_{t=0} &= u_0\end{aligned}$$

We take  $f(s) = s^3 - s$  (more generally :  $f(s) = \sum_{i=1}^{2p+1} a_i s^i, a_{2p+1} > 0$ ),  $h \equiv 0$

First well-posedness result : A. Bertozzi et al.

To go further : global in time/dissipative estimate

First step : obtain an estimate in  $H^{-1}(\Omega)$

→ We need to estimate  $\langle u \rangle = \frac{1}{\text{Vol}(\Omega)} \int_{\Omega} u \, dx$

Classical Cahn-Hilliard equation : conservation of mass

If  $|\langle u_0 \rangle| \leq M$ , then  $|\langle u(t) \rangle| \leq M, t \geq 0$

Equation for  $\langle u \rangle$  :

$$\frac{d\langle u \rangle}{dt} + \frac{1}{\text{Vol}(\Omega)} \int_{\Omega \setminus D} u \, dx = 0$$

We set

$$u = \langle u \rangle + v$$

We find

$$\frac{d\langle u \rangle}{dt} + c_0 \langle u \rangle = -\frac{1}{\text{Vol}(\Omega)} \int_{\Omega \setminus D} v \, dx, \quad c_0 = \frac{\text{Vol}(\Omega \setminus D)}{\text{Vol}(\Omega)}$$

$v$  is solution to

$$\begin{aligned} & \frac{\partial}{\partial t} (-\Delta)^{-1} v - \Delta v + f(\langle u \rangle + v) - \langle f(\langle u \rangle + v) \rangle \\ & + (-\Delta)^{-1} (\chi_{\Omega \setminus D}(x) u - \langle \chi_{\Omega \setminus D}(x) u \rangle) = 0 \end{aligned}$$

$(-\Delta)^{-1}$  : inverse minus Laplacian acting on functions with null average  
( $\langle v \rangle = 0$ )

Multiply the equation by  $v$

Use the inequality

$$\begin{aligned} & ((f(\langle u \rangle + v) - \langle f(\langle u \rangle + v) \rangle, v))_{L^2} \\ &= ((f(\langle u \rangle + v) - f(\langle u \rangle), v))_{L^2} \\ &\geq \frac{c_0}{2} \int_{\Omega} (v^4 + v^2 \langle u \rangle^2) dx - \|v\|_{L^2}^2 \end{aligned}$$

We obtain

$$\frac{d}{dt} \|v\|_{H^{-1}}^2 + \|\nabla v\|_{L^2}^2 + c_0 \int_{\Omega} (v^4 + v^2 \langle u \rangle^2) dx \leq c$$

Consequence :  $\|v\|_{H^{-1}}^2 \leq e^{-ct} \|v_0\|_{H^{-1}}^2 + c', c > 0, t \geq 0$

Multiply the equation by  $-\Delta v$  :

$$\begin{aligned}\|v(t)\|_{L^2} &\leq Q(\|u_0\|_{L^2}), \quad t \geq 0 \\ \|v(t)\|_{L^2} &\leq c, \quad t \geq t_0, \quad t_0 > 0\end{aligned}$$

$Q$  : monotone increasing function

$c$  : independent of  $u_0$  and  $t$ ,  $\|u_0\|_{L^2} \leq R$ ,  $t_0 = t_0(R)$

Equation for  $\langle u \rangle$  :

$$\frac{d\langle u \rangle}{dt} + c_0 \langle u \rangle = -\frac{1}{\text{Vol}(\Omega)} \int_{\Omega \setminus D} v \, dx$$

$$\rightarrow |\langle u(t) \rangle| \leq Q(\|u_0\|_{L^2})e^{-ct} + c', \quad c > 0, \quad t \geq 0$$

Well-posedness, further regularity

Existence of finite-dimensional attractors

Open problem : convergence of solutions to steady states

Numerical simulations : one-step algorithm with threshold

Two-step algorithm :  $\epsilon = 0.1$  and then  $\epsilon = 0.01$

Here :  $\epsilon = 0.05$  and then threshold

If  $u \geq \frac{1}{2}$ , then we take  $u = 1$

If  $u < \frac{1}{2}$ , then we take  $u = 0$

When  $D$  is not "too large" : results comparable with the two-steps algorithm, computation time divided by two







## Logarithmic nonlinear terms :

$$\begin{aligned}\frac{\partial u}{\partial t} + \Delta^2 u - \Delta f(u) + \chi_{\Omega \setminus D}(x)u &= 0 \\ \frac{\partial u}{\partial \nu} &= \frac{\partial \Delta u}{\partial \nu} = 0 \text{ on } \Gamma \\ u|_{t=0} &= u_0\end{aligned}$$

$$f(s) = -2\theta_0 s + \theta_1 \ln \frac{1+s}{1-s}, \quad s \in (-1, 1), \quad 0 < \theta_1 < \theta_0$$

For  $h \neq 0$ , we need  $\int_{\Omega \setminus D} h \, dx = 0$

We have a local (in time) existence result

**Theorem :** We assume that  $u_0 \in H^1(\Omega)$ ,  $|\langle u_0 \rangle| < 1$  and  $-1 < u_0(x) < 1$  a.e.  $x \in \Omega$ . Then, there exists  $T_0 = T_0(u_0)$  and a solution  $u$  such that  $u \in \mathcal{C}([0, T_0]; H^{-1}(\Omega)) \cap L^\infty(0, T_0; H^1(\Omega)) \cap L^2(0, T_0; H^2(\Omega))$  and  $\frac{\partial u}{\partial t} \in L^2(0, T_0; H^{-1}(\Omega))$ . Furthermore,

$$-1 < u(t, x) < 1 \text{ a.e. } (t, x) \in (0, T_0) \times \Omega.$$

We approximate the singular nonlinear term by regular ones

The approximated functions need to satisfy a (uniform) coercivity relation

We set  $F_N(s) = F_{1,N}(s) - \theta_0 s^2$

$$F_{1,N}(s) = \begin{cases} \sum_{k=0}^4 \frac{1}{k!} F_1^{(k)}(1 - \frac{1}{N})(s - 1 + \frac{1}{N})^k, & s \geq 1 - \frac{1}{N}, \\ F_1(s), & |s| \leq 1 - \frac{1}{N}, \\ \sum_{k=0}^4 \frac{1}{k!} F_1^{(k)}(-1 + \frac{1}{N})(s + 1 - \frac{1}{N})^k, & s \leq -1 + \frac{1}{N}. \end{cases}$$

$$F(s) = F_1(s) - \theta_0 s^2$$

$$F_{1,N}(s) = \sum_{k=0}^4 \frac{1}{k!} F_1^{(k)}(1 - \frac{1}{N})(s - 1 + \frac{1}{N})^k, \quad s > 1 - \frac{1}{N}$$

$$F_{1,N}(s) = F_1(s), \quad |s| \leq 1 - \frac{1}{N}$$

$$F_{1,N}(s) = \sum_{k=0}^4 \frac{1}{k!} F_1^{(k)}(-1 + \frac{1}{N})(s + 1 - \frac{1}{N})^k, \quad s < -1 + \frac{1}{N}$$

The approximated functions  $f_N = F'_N$  satisfy

- $F_N \in \mathcal{C}^4(R)$
- $f_N(0) = 0$
- $f'_N \geq -\theta_0, F_N \geq -c_1, c_1 \geq 0$
- $f_N(s)s \geq c_2(F_N(s) + |f_N(s)|) - c_3, c_2 > 0, c_3 \geq 0, s \in R$
- $(f_N(s+a) - f_N(a))s \geq c_4(s^4 + a^2s^2) - c_5, c_4 > 0, c_5 \geq 0, s, a \in R$

All the constants are independent of  $N$

Approximated problems :

$$\begin{aligned}\frac{\partial u_N}{\partial t} + \Delta^2 u_N - \Delta f_N(u_N) + \chi_{\Omega \setminus D}(x)(u_N - h) &= 0 \\ \frac{\partial u_N}{\partial \nu} &= \frac{\partial \Delta u_N}{\partial \nu} = 0 \text{ on } \Gamma \\ u_N|_{t=0} &= u_0\end{aligned}$$

We have the well-posedness and the regularity of the solutions to the approximated problems

A priori estimates :

Equation for the spatial average :

$$\frac{d\langle u_N \rangle}{dt} + \frac{1}{\text{Vol}(\Omega)} \int_{\Omega \setminus D} u_N dx = 0$$

Set  $u_N = \langle u_N \rangle + \bar{u}_N$  :

$$\frac{d\langle u_N \rangle}{dt} + c\langle u_N \rangle = -\frac{1}{\text{Vol}(\Omega)} \int_{\Omega \setminus D} \bar{u}_N dx$$

$$c = \frac{\text{Vol}(\Omega \setminus D)}{\text{Vol}(\Omega)}$$



$\bar{u}_N$  is solution to

$$\frac{\partial \bar{u}_N}{\partial t} + \Delta^2 \bar{u}_N - \Delta(f_N(u_N) - \langle f_N(u_N) \rangle) + \chi_{\Omega \setminus D}(x)u_N - \langle \chi_{\Omega \setminus D}(x)u_N \rangle = 0$$

$$\frac{\partial \bar{u}_N}{\partial \nu} = \frac{\partial \Delta \bar{u}_N}{\partial \nu} = 0 \text{ on } \Gamma$$

$$\bar{u}_N|_{t=0} = v_0 = u_0 - \langle u_0 \rangle$$

Equivalent formulation :

$$\begin{aligned} & (-\Delta)^{-1} \frac{\partial \bar{u}_N}{\partial t} - \Delta \bar{u}_N + f_N(u_N) - \langle f_N(u_N) \rangle \\ & + (-\Delta)^{-1} (\chi_{\Omega \setminus D}(x)u_N - \langle \chi_{\Omega \setminus D}(x)u_N \rangle) = 0 \\ & \frac{\partial \bar{u}_N}{\partial \nu} = 0 \text{ on } \Gamma \end{aligned}$$

Multiply by  $\bar{u}_N$  :

$$\begin{aligned} & \frac{1}{2} \frac{d}{dt} \|\bar{u}_N\|_{-1}^2 + \|\bar{u}_N\|_V^2 \\ & + ((f_N(u_N) - \langle f_N(u_N) \rangle), \bar{u}_N)) + ((\chi_{\Omega \setminus D}(x) u_N, (-\Delta)^{-1} \bar{u}_N)) = 0 \end{aligned}$$

Note that

$$((f_N(u_N) - \langle f_N(u_N) \rangle), \bar{u}_N)) = ((f_N(u_N) - f_N(\langle u_N \rangle), \bar{u}_N))$$

Thus :

$$((f_N(u_N) - \langle f_N(u_N) \rangle), \bar{u}_N)) \geq c_4 (\|\bar{u}_N\|_{L^4(\Omega)}^4 + \langle u_N \rangle^2 \|\bar{u}_N\|^2) - c$$

Furthermore :

$$\begin{aligned} |((\chi_{\Omega \setminus D}(x) u_N, (-\Delta)^{-1} \bar{u}_N))| & \leq c (\|\bar{u}_N\|^2 + |\langle u_N \rangle| \|\bar{u}_N\|) \\ & \leq \frac{c_4}{2} (\|\bar{u}_N\|_{L^4(\Omega)}^4 + \langle u_N \rangle^2 \|\bar{u}_N\|^2) + c \end{aligned}$$

Thus :

$$\frac{d}{dt} \|\bar{u}_N\|_{-1}^2 + \|\bar{u}_N\|_V^2 + c_4(\|\bar{u}_N\|_{L^4(\Omega)}^4 + \langle u_N \rangle^2 \|\bar{u}_N\|^2) \leq c$$

Next :

$$\frac{d\langle u_N \rangle^2}{dt} + c_6 \langle u_N \rangle^2 \leq c \|\bar{u}_N\|^2$$

Thus :

$$\frac{d\langle u_N \rangle^2}{dt} + c_6 \langle u_N \rangle^2 \leq \frac{c_4}{2} (\|\bar{u}_N\|_{L^4(\Omega)}^4 + \langle u_N \rangle^2 \|\bar{u}_N\|^2) + c$$

Sum :

$$\frac{dE_{1,N}}{dt} + c(\|u_N\|_{H^1(\Omega)}^2 + \|\bar{u}_N\|_{L^4(\Omega)}^4 + \langle u_N \rangle^2 \|\bar{u}_N\|^2) \leq c', \quad c > 0$$

$$E_{1,N} = \langle u_N \rangle^2 + \|\bar{u}_N\|_{-1}^2$$

$$E_{1,N} \geq c \|u_N\|_{H^{-1}(\Omega)}^2, \quad c > 0$$

Multiply the original equation by  $u_N$  :

$$\frac{d}{dt} \|u_N\|^2 + \|\Delta u_N\|^2 \leq 2c_0 \|\nabla u_N\|^2 + c \|u_N\|^2$$

Combine the two estimates :

$$\frac{dE_{2,N}}{dt} + c(\|u_N\|_{H^2(\Omega)}^2 + \|\bar{u}_N\|_{L^4(\Omega)}^4 + \langle u_N \rangle^2 \|\bar{u}_N\|^2) \leq c', \quad c > 0$$

$$E_{2,N} = \delta_1 \|u_N\|^2 + E_{1,N}$$

$$E_{2,N} \geq c \|u_N\|^2, \quad c > 0$$

Equivalent equations :

$$\frac{\partial u_N}{\partial t} + \chi_{\Omega \setminus D}(x) u_N = \Delta \mu_N$$

$$\mu_N = -\Delta u_N + f_N(u_N)$$

$$\frac{\partial u_N}{\partial \nu} = \frac{\partial \mu_N}{\partial \nu} = 0 \text{ on } \Gamma$$

Multiply the first equation by  $\mu_N$  and the second by  $\frac{\partial u_N}{\partial t}$  :

$$\frac{1}{2} \frac{d}{dt} (\|\nabla u_N\|^2 + 2 \int_{\Omega} F_N(u_N) dx) + \|\nabla \mu_N\|^2 = -((u_N, \chi_{\Omega \setminus D}(x) \mu_N))$$

Multiply the second equation by  $\chi_{\Omega \setminus D}(x) u_N$  :

$$((u_N, \chi_{\Omega \setminus D}(x) \mu_N)) = -((\Delta u_N, \chi_{\Omega \setminus D}(x) u_N)) + \int_{\Omega \setminus D} f_N(u_N) u_N dx$$

Thus :

$$\begin{aligned} & \frac{d}{dt} (\|\nabla u_N\|^2 + 2 \int_{\Omega} F_N(u_N) dx) \\ & + c (\|\nabla \mu_N\|^2 + \int_{\Omega \setminus D} |f_N(u_N)| dx + \int_{\Omega \setminus D} F_N(u_N) dx) \leq c' \|u_N\|_{H^2(\Omega)}^2 + c'', \quad c > 0 \end{aligned}$$

Combining :

$$\begin{aligned} & \frac{dE_{3,N}}{dt} + c(\|u_N\|_{H^2(\Omega)}^2 + \|\bar{u}_N\|_{L^4(\Omega)}^4 + \langle u_N \rangle^2 \|\bar{u}_N\|^2 \\ & + \int_{\Omega \setminus D} |f_N(u_N)| dx + \int_{\Omega \setminus D} F_N(u_N) dx + \|\nabla \mu_N\|^2) \leq c', \quad c > 0 \end{aligned}$$

$$E_{3,N} = \delta_2(\|\nabla u_N\|^2 + 2 \int_{\Omega} F_N(u_N) dx) + E_{2,N}$$

$$E_{3,N} \geq c\|u_N\|_{H^1(\Omega)}^2 - c', \quad c > 0$$

Equivalent equations :

$$(-\Delta)^{-1} \frac{\partial \bar{u}_N}{\partial t} + (-\Delta)^{-1} (\chi_{\Omega \setminus D}(x) u_N - \langle \chi_{\Omega \setminus D}(x) u_N \rangle) = -(\mu_N - \langle \mu_N \rangle)$$

$$\mu_N - \langle \mu_N \rangle = -\Delta \bar{u}_N + f_N(u_N) - \langle f_N(u_N) \rangle$$

Thus :

$$\left\| \frac{\partial \bar{u}_N}{\partial t} \right\|_{-1} \leq c(\|u_N\| + \|\nabla \mu_N\|)$$

and

$$\left\| \frac{\partial u_N}{\partial t} \right\|_{H^{-1}(\Omega)} \leq c(\|u_N\| + \|\nabla \mu_N\|)$$

Furthermore :

$$\|f_N(u_N) - \langle f_N(u_N) \rangle\| \leq c(\|u_N\|_{H^2(\Omega)} + \|\nabla \mu_N\|)$$

Finally :

$$\begin{aligned} \frac{dE_{3,N}}{dt} + c(\|u_N\|_{H^2(\Omega)}^2 + \|\bar{u}_N\|_{L^4(\Omega)}^4 + \langle u_N \rangle^2 + \|\bar{u}_N\|^2 \\ + \|\frac{\partial u_N}{\partial t}\|_{-1}^2 + \|f_N(u_N) - \langle f_N(u_N) \rangle\|^2 \\ + \int_{\Omega \setminus D} |f_N(u_N)| dx + \int_{\Omega \setminus D} F_N(u_N) dx) \leq c', \quad c > 0 \end{aligned}$$

$$\bar{u}_N = u_N - \langle u_N \rangle$$

$$\begin{aligned} E_{3,N} &= \|u_N\|_{-1}^2 + \delta_1 \|u_N\|^2 + \delta_2 (\|\nabla u_N\|^2 + 2 \int_{\Omega} F_N(u_N) dx) \\ \delta_1, \delta_2 &> 0 \text{ small} \end{aligned}$$

$$E_{3,N} \geq c \|u_N\|_{H^1(\Omega)}^2 - c', \quad c > 0$$



Not sufficient to pass to the limit

We need to estimate  $|\langle f_N(u_N) \rangle|$  (hence an estimate on the  $L^2$ -norm of  $f_N(u_N)$ )

To do so, we need an estimate of the form  $|\langle u_N \rangle| \leq 1 - \delta$ ,  $\delta \in (0, 1)$   
independent of  $N$

We can prove this only locally in time

We assume that  $|\langle u_0 \rangle| \leq 1 - 2\delta$ ,  $\delta > 0$  given

We have

$$\frac{d\langle u_N \rangle}{dt} + c\langle u_N \rangle = -\frac{1}{\text{Vol}(\Omega)} \int_{\Omega \setminus D} \bar{u}_N dx$$

This yields

$$\langle u_N(t) \rangle = e^{-ct} \langle u_0 \rangle - e^{-ct} \int_0^t e^{cs} ds \int_{\Omega \setminus D} \bar{u}_N dx$$

and

$$\begin{aligned} |\langle u_N(t) \rangle| &\leq |\langle u_0 \rangle| + ce^{-ct} \int_0^t e^{cs} \|u_N\| ds \\ &\leq 1 - 2\delta + c'(1 - e^{-ct}) \end{aligned}$$

→ There exists  $T_0 = T_0(u_0, \delta) > 0$  independent of  $N$  such that

$$|\langle u_N(t) \rangle| \leq 1 - \delta, \quad t \in [0, T_0]$$

We can then prove that

$$|\langle f_N(u_N) \rangle| \leq c_\delta \|\bar{u}_N\|_{L^2} \|f_N(u_N) - \langle f_N(u_N) \rangle\|_{L^2} + c'_\delta$$

→ Uniform estimate on the  $L^2$ -norm of  $f_N(u_N)$  on  $[0, T_0]$

This allows to pass to the limit on  $[0, T_0]$

Existence of a local (in time) solution

**Remark :** We rewrite the equation in the form

$$\frac{\partial u}{\partial t} + \Delta^2 u - \Delta f(u) + u - \chi_D(x)u = 0$$

We then have

$$\frac{d\langle u \rangle}{dt} + \langle u \rangle = \frac{1}{\text{Vol}(\Omega)} \int_D u \, dx$$

and

$$|\langle u(t) \rangle| \leq e^{-t} |\langle u_0 \rangle| + \frac{\text{Vol}(D)}{\text{Vol}(\Omega)} e^{-t} \int_0^t e^s \, ds$$

This yields

$$|\langle u(t) \rangle| \leq \max(|\langle u_0 \rangle|, \frac{\text{Vol}(D)}{\text{Vol}(\Omega)})$$

whence

$$|\langle u(t) \rangle| \leq 1 - \delta$$

where  $\delta = \delta(u_0) \in (0, 1)$  is independent of time

→ The solutions are global in time

Numerical simulations :

One-step algorithm with threshold

The convergence time is faster

The results are better with large inpainting domains



## Extensions of the model :

**Cahn-Hilliard inpainting for multicolor images** (L. Cherfils-H. Fakh-A. Miranville, JMIV) :

Multiphasic Cahn-Hilliard system

Each phase corresponds to a color

We consider the hyperplane

$$S = \{c \in R^n \text{ such that } \sum_{i=1}^n c_i = 1\}$$

$h = (h_1, \dots, h_n) \in S$  : damaged image, known on  $\Omega \setminus D$



We look for  $u = (u_1, \dots, u_n) \in S$  such that

$$\begin{aligned}\frac{\partial u_i}{\partial t} &= \Delta \mu_i + \lambda_0 \chi_{\Omega \setminus D}(x)(h_i - u_i), \quad i = 1, \dots, n \\ \mu_i &= f_i(u) - \epsilon^2 \Delta u_i, \quad i = 1, \dots, n \\ \frac{\partial u_i}{\partial \nu} &= \frac{\partial \mu_i}{\partial \nu} = 0 \text{ on } \Gamma, \quad i = 1, \dots, n \\ u_i|_{t=0} &= u_{i,0}, \quad i = 1, \dots, n\end{aligned}$$

$$f_i(u) = \frac{\partial F(u)}{\partial u_i} - \frac{1}{n} \sum_{j=1}^n \frac{\partial F(u)}{\partial u_j}, \quad i = 1, \dots, n$$

$$F(u) = \frac{1}{n} \sum_{i=1}^n u_i^2 (1 - u_i^2)$$

Lagrange multiplier to ensure  $u \in S$

Well-posedness and regularity of the solutions

Existence of finite-dimensional attractors

The model is algebraically consistent with the diphasic model

Numerical simulations :

One-step algorithm with threshold

Drawback : not efficient when the number of colors  $n$  is large



## **Grayscale Cahn-Hilliard inpainting** : (L. Cherfils-H. Fakh-A. Miranville, MMS)

Aim : propose a simple model

Known models : heavy to implement numerically

J. Bosch et al. : multiphase Cahn-Hilliard system ;  $n$  : number of shades of gray (not efficient when  $n$  is large)

Other models : total variation in  $H^{-1}$ , Low Curvature Image Simplifier (similar drawbacks)

Idea : consider a complex version of the Bertozzi et al. model (H. Grossauer-O. Sherzer : complex Allen-Cahn equation)

$h_1 \in L^2(\Omega)$  : damaged image ( $h_1 : \Omega \rightarrow [-1, 1]$ )

We introduce  $h : \Omega \rightarrow \mathbb{C}$  defined by

$$h = h_1 + ih_2, \quad h_2(x) = \sqrt{1 - h_1(x)^2}$$

$\rightarrow h \in L^2(\Omega; \mathbb{C}), |h| = 1$

Complex version of the Bertozzi et al. model :

$$\begin{aligned} \frac{\partial u}{\partial t} + \epsilon \Delta^2 u - \frac{1}{\epsilon} \Delta f(u) + \lambda_0 \chi_{\Omega \setminus D}(x)(u - h) &= 0 \\ \frac{\partial u}{\partial \nu} &= \frac{\partial \Delta u}{\partial \nu} = 0 \text{ on } \Gamma \\ u|_{t=0} &= u_0 \end{aligned}$$

$$f(z) = |z|^2 z - z, \quad z \in \mathbb{C}$$

Well-posedness and regularity of the solutions

Existence of finite-dimensional attractors

Numerical simulations :

Two-steps algorithm

We use the information on the image known outside the inpainting region

Inpainting result : real part of the solution

We only need to compute two functions whatever the number of shades of gray is











